

**STATE DENSITY FORMALISM OF
THE IWAMOTO-HARADA MODEL:
A SUITABLE TOOL TO TREAT
CLUSTER EMISSION FROM
HEAVY-ION COLLISIONS WITH
ACCOUNT FOR SPIN VARIABLES**

E. Běták

Institute of Physics SAS, Bratislava, Slovakia

J. Cseh

MTA ATOMKI, Debrecen, Hungary

Exotic nuclear shapes (super and hyperdeformed) expected also for some (light) 4- α nuclei, like ^{36}Ar etc. Their formation possible (with probably higher cross section – higher angular momenta) in heavy ion induced reactions, obviously dependent on the way of production. This is a practical stimulation for our effort.

Iwamoto & Harada (1982); Betak & Dobes (1977)

classical coalescence

$$\lambda_c(x, n, \varepsilon) \propto \gamma_x \frac{\omega(p - p_x, h, U)}{\omega(p, h, E)} \frac{\omega(p_x, 0, \varepsilon_x + B_x)}{g_x}$$

IH coalescence (pickup + knockout + coalescence)

$$\sum_{p^*=1}^{p_x} \omega(p - p^*, h, E - \varepsilon_1) \omega(p^*, 0, \varepsilon_1) \omega(0, p_x - p^*, \varepsilon_2) d\varepsilon_1$$

Heavy ions “translated” into the language of the pre-equilibrium exciton model mean replacing the initial exciton number by another value.

Similarly as in the beginning of the exciton (systematics of n_0), we use similar approach for heavy ions. Both very rough derivation and analyses of data (Cindro *et al.*, Ma *et al.*) give similar results.

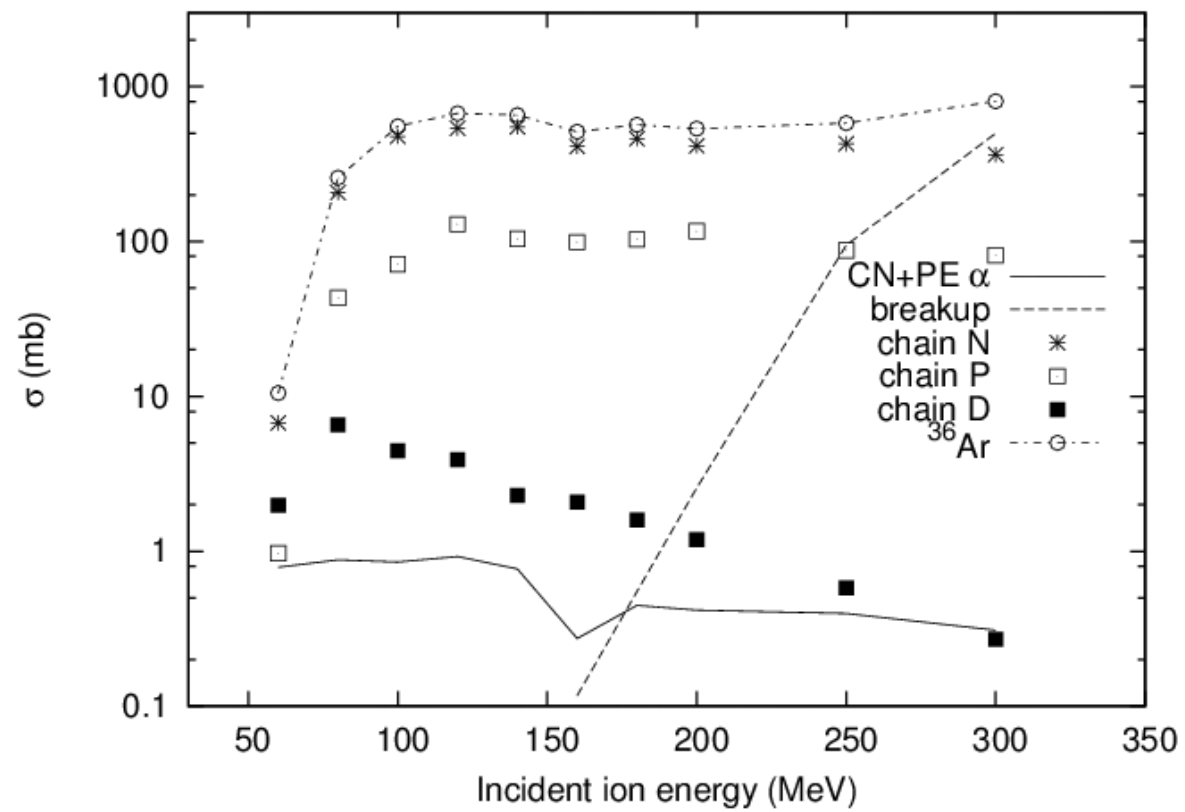
For low energies

$$\frac{n_0}{A_{proj}} = 0.09 + (0.38 - 0.08 \frac{A_{targ} - A_{proj}}{A_{targ} + A_{proj}}) \sqrt{\frac{E_{cm} - V_C}{A_{proj}}}$$

For higher ones

$$\frac{n_0}{E_{exc}} = 4.6 + 0.54 \sqrt{\frac{E_{cm} - V_C}{A_{proj}}}$$

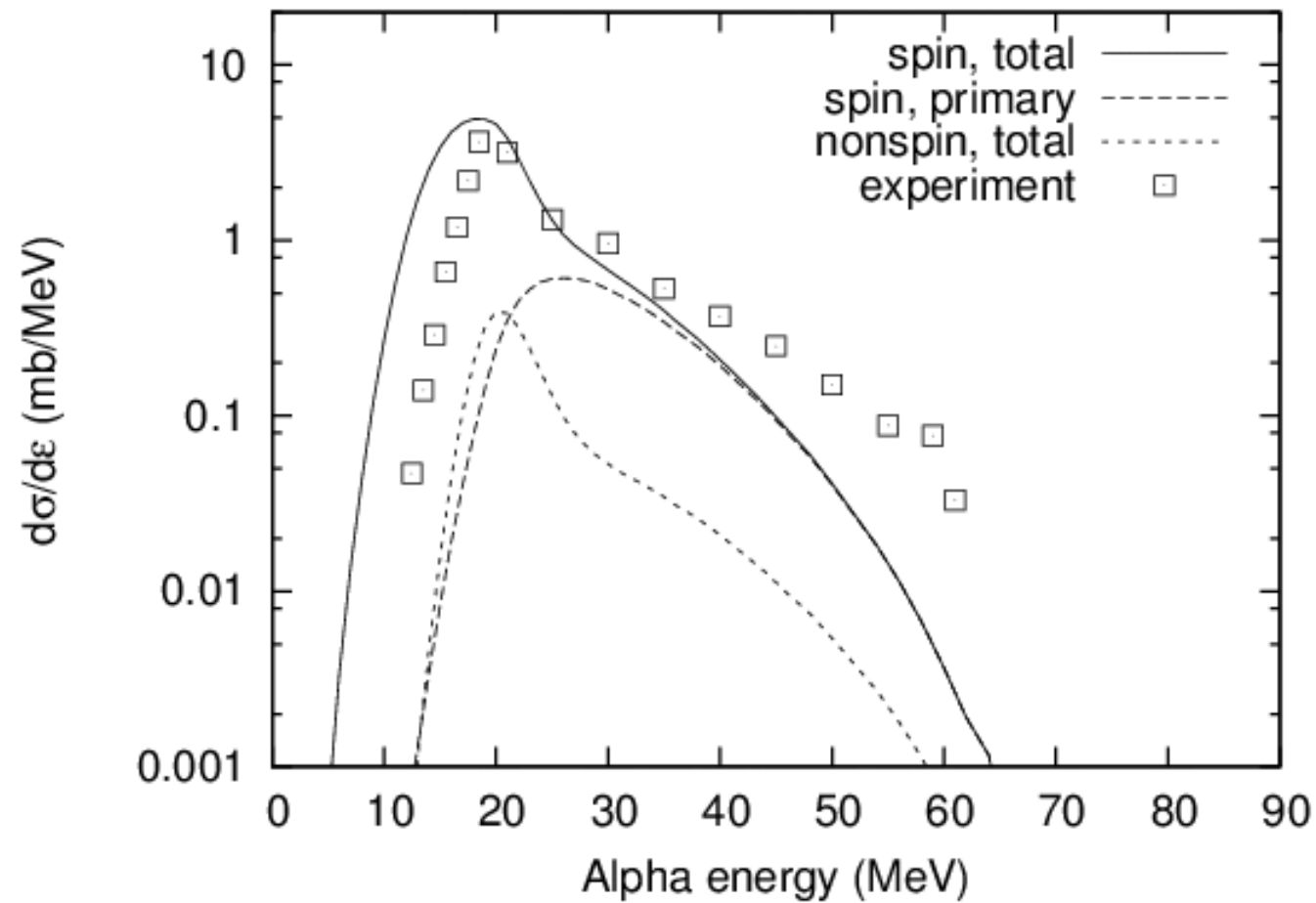
(or replacing 4.6 by 6.6 for “pure systematics”)

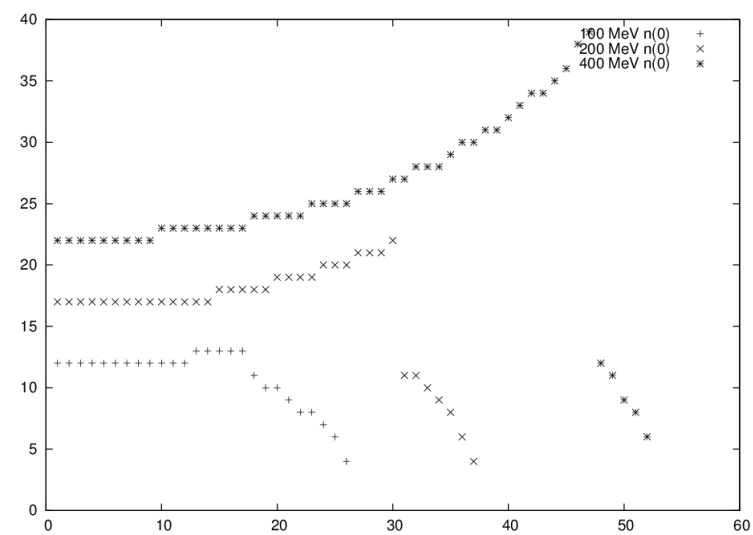
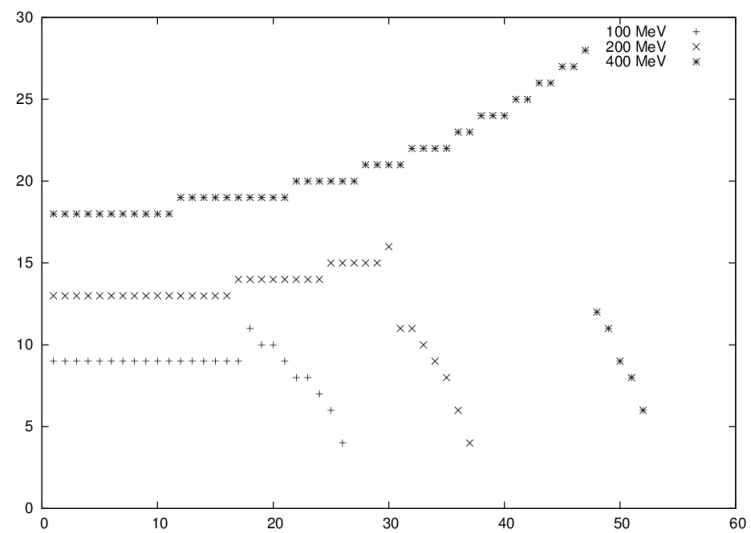


Inclusion of spin by deducting rotational energy from the excitation one

$$E_{exc}(l) = E_{cm} + B - E_{rot}(l)$$

where $E_{rot}(l)$ is calculated with inertia moment of two touching spheres, at least for incident energies ca 5 – 20 MeV/A (*see time-dependent calculations of initial phase of heavy-ion collisions*)





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**Thank you
for your attention**

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